

ME 115a - Lecture Notes

Agenda

★ Office Hour Scheduling

★ Recap of Recent Stuff

★ New Stuff

∞ Twist Coordinates

∞ Rigid Body Velocity

□ Rotational

□ Translational

□ as "Twists"

∞ "Types" of Velocity

□ Body Velocity

□ Spatial Velocity

Recap

→ A fancy french fellow named Charles showed that any spatial displacement of a Rigid body was equivalent to a 'screw motion',

So called because of the similarity to the

motion of a screw ()

Parameters: $(\vec{\omega}, \phi, d'', \vec{p})$
 axis $\vec{\omega}$, rotation angle ϕ , translation amount d'' , axis base position $\vec{p}$

Conversion into homogeneous transformation:

$$g = \begin{bmatrix} e^{\hat{\omega}\phi} & d''\vec{\omega} + (I - A)\vec{s}_\perp \\ \vec{0}^T & 1 \end{bmatrix}$$

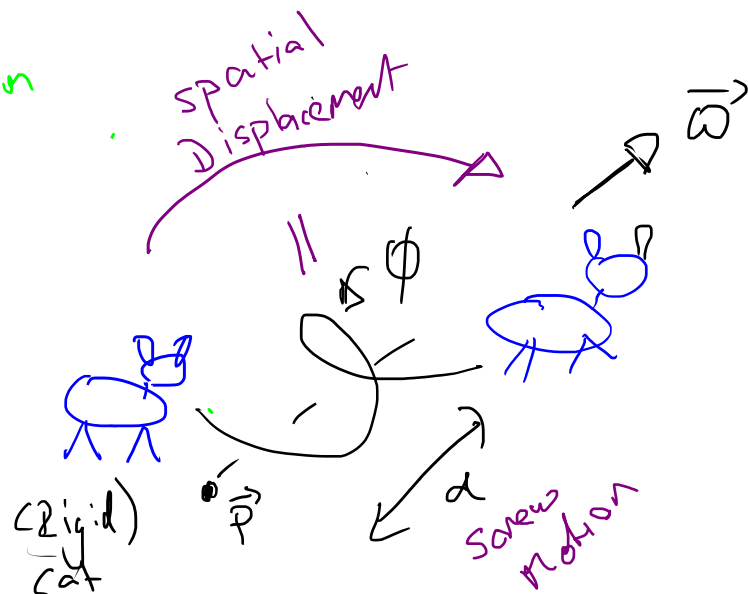


Fig 1. A screwy rigid cat

New Stuff

Twist Coord. for Spatial Displacements

Let $g = \begin{bmatrix} R & d \\ \vec{0}^T & 1 \end{bmatrix}$ be a homogeneous transformation.
 notation: R (rotation), d (translation)

Suppose that for some vector

$$\xi = \begin{bmatrix} \vec{v} \\ \vec{\omega} \end{bmatrix} \in \mathbb{R}^6, \text{ \& some } \phi \in \mathbb{R}, g = \exp(\hat{\xi}\phi)$$

Recall/note that $\hat{\xi} = \begin{bmatrix} \hat{\omega} & \vec{v} \\ \vec{0}^T & 1 \end{bmatrix}$

(calculating $\exp(\phi\hat{\xi})$), we get $\exp(\phi\hat{\xi}) = \begin{bmatrix} e^{\phi\omega} & (I - e^{\phi\hat{\omega}})(\vec{\omega} \times \vec{v}) + \vec{\omega}\vec{\omega}^T \vec{v}\phi \\ \vec{0}^T & 1 \end{bmatrix}$

equating to g in terms of screw parameters,
we get

$$\begin{bmatrix} e^{\hat{\omega}\phi} & | & d'\bar{\omega} + (I - e^{\hat{\omega}\phi})\bar{\omega} \\ \hline \bar{o}' & | & 1 \end{bmatrix} = \begin{bmatrix} e^{\phi\omega} & | & (I - e^{\phi\hat{\omega}})(\bar{\omega}\bar{\omega}) + \bar{\omega}\bar{\omega}^T v \phi \\ \hline \bar{o}'^T & | & 1 \end{bmatrix}$$

If you guessed that $\bar{v}'$ takes the form

$$\bar{v}' = \bar{v}' \times \bar{\omega}' + h\bar{\omega}', \quad (\bar{\omega}' = \bar{\omega})$$

you (would be weird, and you would) find that equality

above holds for $\bar{v}' = \bar{v}'_{\perp} + \underbrace{\langle \bar{\omega} \rangle}_{\in \mathbb{R}, \text{ arbitrary.}}$

SANITY CHECK

Why's the exponential showing up?

Well... exponentials show up all the time, but

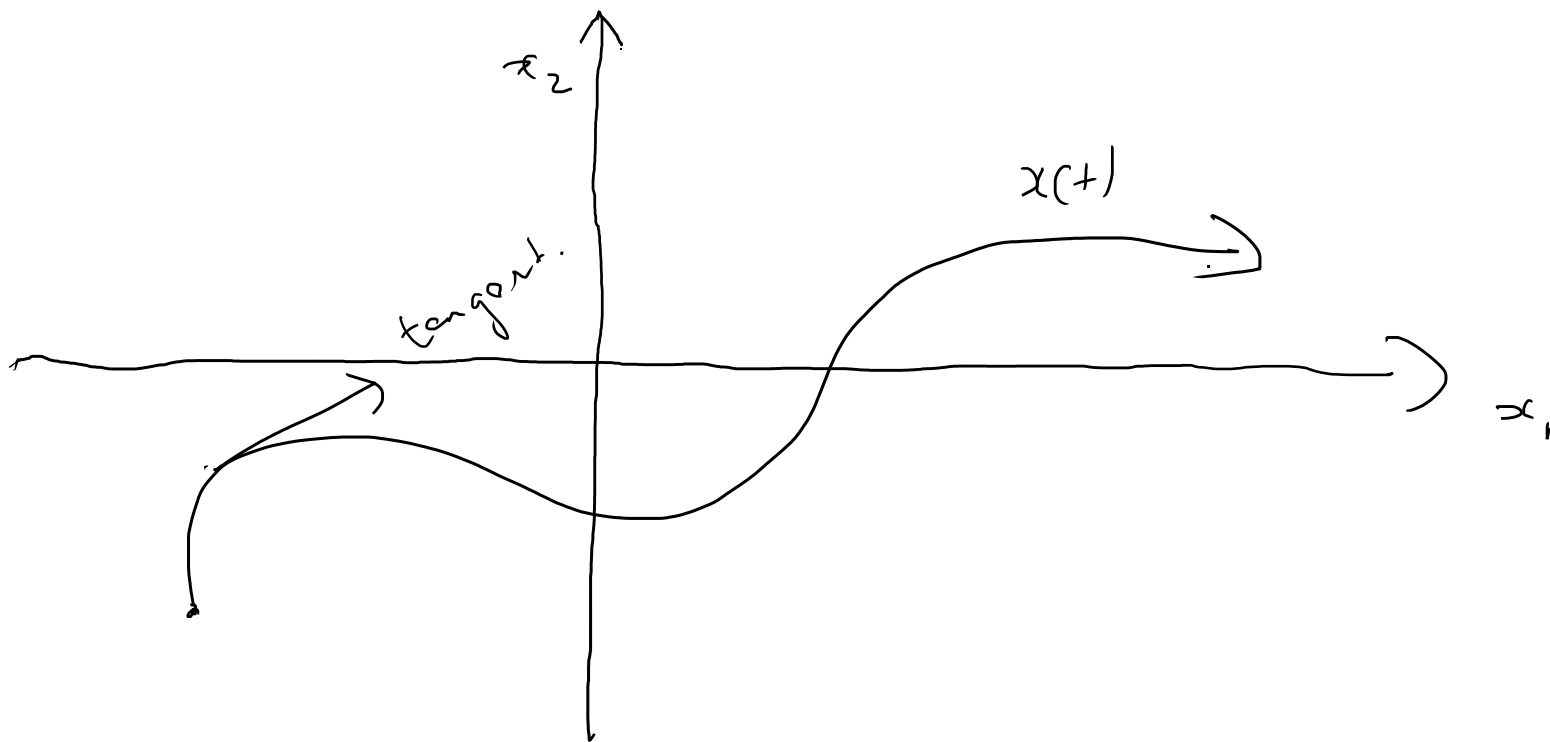
in particular, consider

$$\underbrace{\dot{\bar{x}}}_{\text{Linear}} = \overset{\text{constant matrix}}{A} \bar{x}$$

(3)

solution is $x(t) = e^{At} x_0$.

suppose that $\bar{x}(t)$ is a position in $\mathbb{R}^2$.



The constant ' A ' describes target vectors! The "derivative" of $\bar{x}$.

now, suppose that $x(t)$ was the path of a rigid body, & included position & orientation as a function of time. We could think of $\bar{x}$ as a homogeneous transformation; then, $\frac{d}{dt}(x)$ would look like / have something to do with $\begin{matrix} \uparrow \\ \text{twist} \\ \downarrow \end{matrix}$!

And on that note . . .

▲ Rigid Body Velocity

"TWIST MATRIX"

Take my word that all this formal construction for spatial displacements is extremely useful.

But something is missing . . . real things don't undergo non-zero displacements instantaneously!

2/9/14

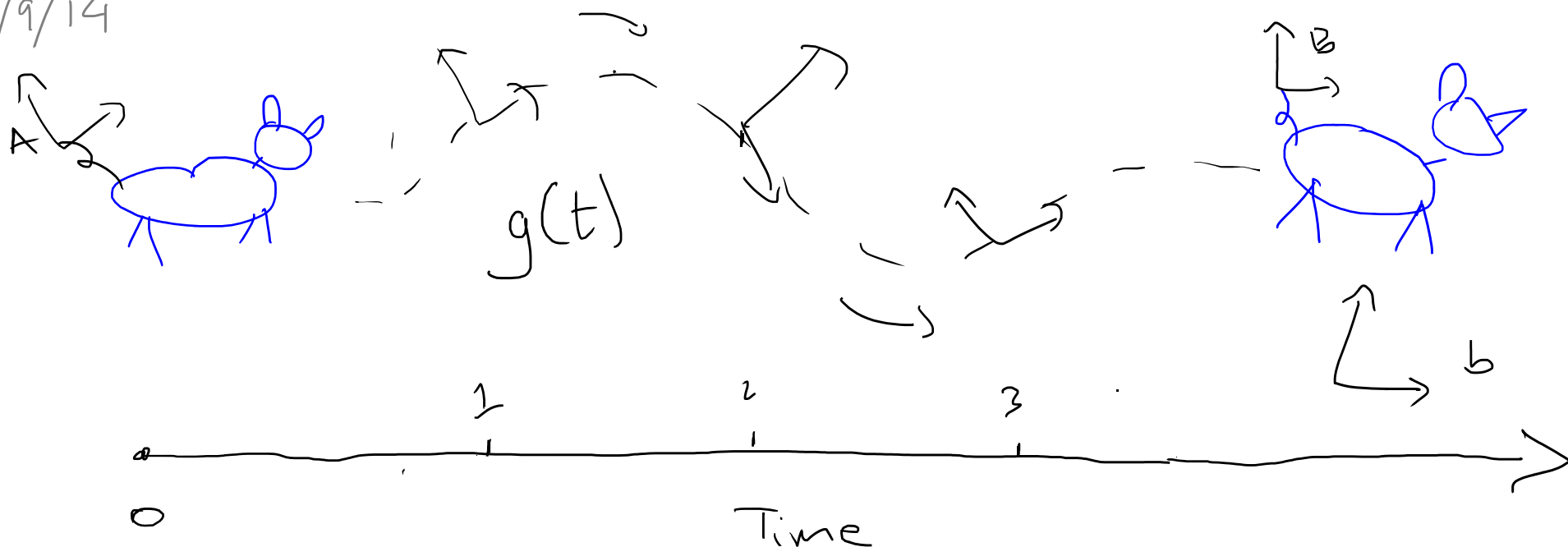
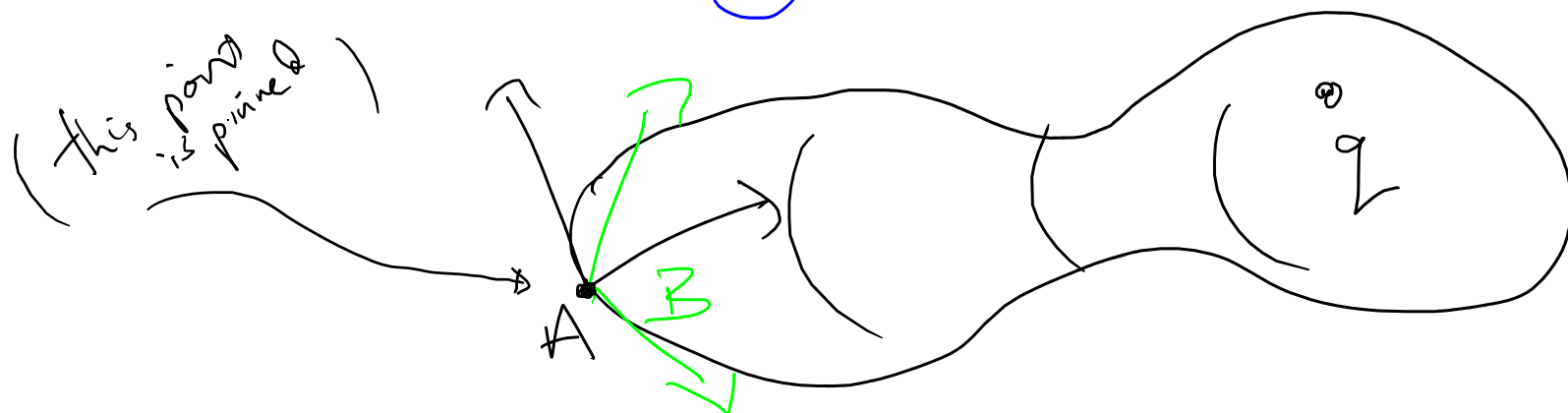


Fig 2. flying (rigid) cat.

Suppose $g(t) = (R(t), p(t))$. Is there a notion of velocity? something like $\dot{g}$?

$\frac{d}{dt}(g(t))$

Rotational Velocity



Suppose a rigid body is undergoing rotation about a fixed point at the origin of a fixed frame, A .

Let B share its origin with A , but rotate along with the body. B is known generally as a "Body frame".

(5)

Let $\underline{R_{ab}}(t) : \mathbb{R}_{\geq 0} \rightarrow SE(3)$ describe the orientation of frame B with respect to frame A at each time t .

For $\underline{q}_b$ a point on the body expressed in the B frame,

point q expressed in terms of frame A

$$\underline{q}_a(t) = R_{ab}(t) \underline{q}_b(t) = R_{ab}(t) \underline{q}_b$$

$\underbrace{\quad}_{\text{constant}}$

lets try to work out velocity:

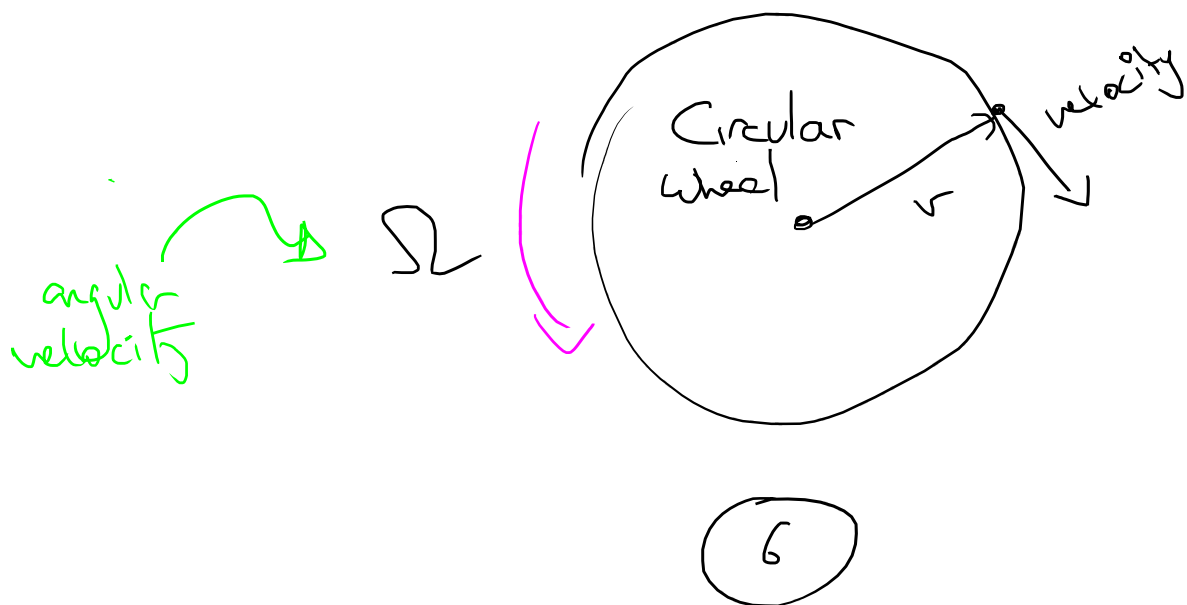
→ in spatial frame: $\underline{V}_{q_a} = \frac{d}{dt} \underline{q}_a(t) = \dot{R}_{ab}(t) \underline{q}_b$

* General note: $[\cdot]_{ab}$ eats things from frame b on the right & spits out things in frame a on the left

map from body coordinates to spatial velocity.

I hope

^ You've seen angular velocity in physics/classical mechanics e.g.



velocity at radius $r = \Omega r$

lets see if we can figure out something
analogous in the case of 3 dimensional rotational velocity:

Recall : $R_{ab}(t) R_{ab}^T(t) = I$ (since $R_{ab} \in SE(3)$)

differentiate
using product rule

$$\dot{R}_{ab} R_{ab}^T + R_{ab} \dot{R}_{ab}^T = 0 \leftarrow (\text{zero matrix})$$

$$\Rightarrow \dot{R}_{ab} R_{ab}^T = - R_{ab} \dot{R}_{ab}^T$$

$$\Rightarrow \dot{R}_{ab} R_{ab}^T \text{ is skew symmetric!}$$

lets use this fact:

$$\begin{aligned} V_{\dot{q}_a}(t) &= \dot{R}_{ab}(t) q_b \\ &= \underbrace{\dot{R}_{ab}(t) R_{ab}^T(t)}_{\text{skew sym bro!}} R_{ab}(t) q_b \end{aligned}$$

lets Define $\hat{\omega}_{ab}^s(t) = \dot{R}_{ab}(t) R_{ab}^T(t)$

Then $V_{\dot{q}_a}(t) = \hat{\omega}_{ab}^s(t) \underbrace{R_{ab}(t) q_b}_{\leftarrow q_a!}$

$$\begin{aligned} V_{\dot{q}_a}(t) &= \hat{\omega}_{ab}^s(t) q_a \\ V_{\dot{q}_a}(t) &= \vec{\omega}_{ab}^s(t) \times q_a \end{aligned}$$

$\vec{\omega}_{ab}^s$ here is the body's

Spatial Angular Velocity

The quantity $R_{ab}^T \dot{R}_{ab}$ is skew symm
by a similar argument to the one above,

Setting $\hat{\omega}_{ab}^b = R_{ab}^T \dot{R}_{ab}$

$$\begin{aligned} V_{q_b}(t) &= R_{ab}^T(t) V_{q_a}(t) \\ &= R_{ab}^T(t) \dot{R}_{ab}(t) q_b \end{aligned}$$

$$\begin{aligned} V_{q_b}(t) &= \hat{\omega}_{ab}^b q_b \\ &= \vec{\omega}_{ab}^b \times q_b \end{aligned}$$

$\vec{\omega}_{ab}^b$ is called the body's **body angular velocity**.

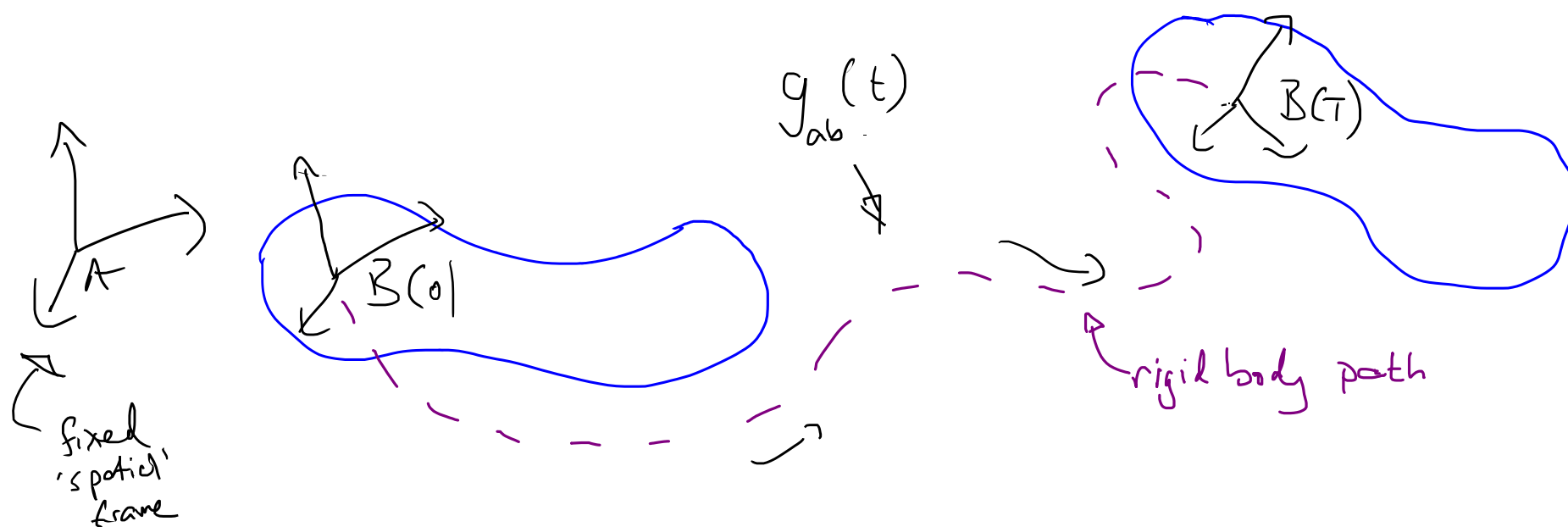
$V_{q_b}(t)$ is the velocity of point q **as seen**

in the B frame

Spatial = (Rotational + Translational) Velocities

It's "easy" to imagine that we might generalise this kind of thing (quantities that capture instantaneous velocity at time t) to arbitrary displacements.

More general picture:



Suppose that

$$g_{ab}(t) = \begin{bmatrix} R_{ab}(t) & d_{ab}(t) \\ 0^T & 1 \end{bmatrix}$$

Let's tinker with derivatives along the same lines as we did for rotations

$$\dot{g}_{ab} g_{ab}^{-1} = \begin{bmatrix} \dot{R}_{ab} & \dot{d}_{ab} \\ 0^T & 0 \end{bmatrix} \begin{bmatrix} R_{ab}^T & -R_{ab}^T d_{ab} \\ 0 & 1 \end{bmatrix}$$

$\dot{g}_{ab}$ g_{ab}^{-1}

$$= \begin{bmatrix} \dot{R}_{ab} R_{ab}^T & (\dot{d}_{ab} - R_{ab} R_{ab}^T d_{ab}) \\ 0^T & 0 \end{bmatrix}$$

Look familiar?

2/9/14

Let's Define

$$\hat{V}_{ab}^s = \begin{bmatrix} \dot{R}_{ab} R_{ab}^T & (\dot{d}_{ab} - R_{ab} R_{ab}^T d_{ab}) \\ 0^T & 0 \end{bmatrix}$$

The "Spatial Velocity" of the Rigid Body.

We can 'un-hat' $\hat{V}_{ab}^s$ to get

$$\bar{V}_{ab}^s = \begin{bmatrix} v_{ab} \\ \omega_{ab} \end{bmatrix} \in \mathfrak{se}(3) \quad \text{where } v_{ab} = (\dot{d}_{ab} - R_{ab} R_{ab}^T d_{ab})$$

$$\hat{\omega}_{ab} = \dot{R}_{ab} R_{ab}^T$$

Remember this?

$\hat{V}_{ab}^s$ is a twist matrix and V_{ab}^s is the

corresponding twist vector. The set of twist

vectors is isomorphic to twist matrices. We can

think of R.B. velocities either way, or call them "twists" to save oxygen.

isomorphism: A bijective map that preserves some structure.

In this case: $\bar{V}_{ab}^s = \underbrace{(\hat{V}_{ab}^s)^V}_{\text{unmap}} = \begin{bmatrix} v_{ab} \\ \omega_{ab} \end{bmatrix}$

The structure that is preserved is the vector space structure of $se(3)$.

How do we interpret spatial velocities?

▣ $\bar{\omega}_{ab}^s$ = instantaneous axis of rotation -

weird Alert! → ▣ $\bar{v}_{ab}^s$ = velocity of particle coincident to origin of frame A -

A little weird! Mull this over a few times!

We can do the same thing we did for rotational motion to get Body Velocities,

Set $\hat{V}_{ab}^b = g_{ab}^{-1} \dot{g}_{ab} = \begin{bmatrix} R_{ab}^T \dot{R}_{ab} & R_{ab}^T \dot{d}_{ab} \\ 0^T & 0 \end{bmatrix}$

$$\hat{V}_{ab}^b = \begin{bmatrix} \hat{\omega}_{ab}^b & \hat{v}_{ab}^b \\ 0^T & 0 \end{bmatrix}$$

OR $\bar{V}_{ab}^b = \left(\hat{V}_{ab}^b \right)^V = \begin{bmatrix} \bar{v}_{ab}^b \\ \bar{\omega}_{ab}^b \end{bmatrix} = \begin{bmatrix} R_{ab}^T \dot{d}_{ab} \\ \bar{\omega}_{ab}^b \end{bmatrix}$

Body velocity: velocity of body as seen in body's own frame.